ECE 340 Lectures 34 Semiconductor Electronics

Spring 2022 10:00-10:50am Professor Umberto Ravaioli Department of Electrical and Computer Engineering 2062 ECE Building

Today's Discussion

- MOS Capacitor (conclusion)
 - Capacitance-Voltage Analysis
 - Realistic effects

We have studied the Ideal MOSFET Capacitor



gate contact oxide

Silicon bulk

contact

Summary of conditions – surface potential



We do not know the exact behavior of $\phi(x)$ but we know the relationship between charge and potential

Poisson equation

$$\frac{d^2\phi}{dx^2} = -\frac{\rho(x)}{\epsilon_s}$$

with charge density

$$\rho(x) = q[N_D^+ - N_A^- + p(x) - n(x)]$$

Electric Field
$$\mathcal{E} = -\frac{d\phi}{dx}$$

From Last Lecture: Analytical model for n(x)

away from interface

$$n_o = n_i \exp\left(\frac{E_F - E_i}{k_B T}\right) = n_i \exp\left(-\frac{q\phi_F}{k_B T}\right)$$

at any x location

$$n(x) = n_i \exp\left(\frac{E_F - E_i(x)}{k_B T}\right) = n_i \exp\left(-q \frac{\phi_F - \phi(x)}{k_B T}\right) =$$
$$= n_i \exp\left(-q \frac{\phi_F}{k_B T}\right) \exp\left(q \frac{\phi(x)}{k_B T}\right) = n_0 \exp\left(q \frac{\phi(x)}{k_B T}\right)$$

From Last Lecture: Analytical model for p(x)

away from interface

$$p_o = n_i \exp\left(\frac{E_i - E_F}{k_B T}\right) = n_i \exp\left(\frac{q\phi_F}{k_B T}\right)$$

at any x location

$$p(x) = n_i \exp\left(\frac{E_i(x) - E_F}{k_B T}\right) = n_i \exp\left(-q \frac{\phi(x) - \phi_F}{k_B T}\right) =$$
$$= n_i \exp\left(q \frac{\phi_F}{k_B T}\right) \exp\left(-q \frac{\phi(x)}{k_B T}\right) = p_0 \exp\left(-q \frac{\phi(x)}{k_B T}\right)$$



$$\frac{d^2\phi}{dx^2} = \frac{d}{dx} \left(\frac{d\phi}{dx} \right) = -\frac{\rho(x)}{\epsilon_s}$$



$$\frac{d^2 \phi}{dx^2} = \frac{d}{dx} \left(\frac{d \phi}{dx} \right) =$$

$$\frac{RHS}{\epsilon_s} \left(\text{Right Hand Side} \right)$$

$$= -\frac{q}{\epsilon_s} \left\{ p_o \left[\exp \left(-\frac{q \phi}{k_B T} \right) \left(-1 \right) \right] - n_o \left[\exp \left(\frac{q \phi}{k_B T} \right) \left(-1 \right) \right] \right\}$$

$$-p_o \left[n_o \left[N_D^+ - N_A^- = n_o - p_o \right] \right]$$

 $\chi \to \infty$

Integrate from the bottom to x = 0any coordinate x, toward the oxide interface, where we have

 $\mathcal{E}(x)$



Solution result:

$$\mathcal{E}^{2} = \frac{2k_{B}T}{\varepsilon_{s}} p_{0} \left[\left(\exp\left(-\frac{q\phi}{k_{B}T}\right) + \frac{q\phi}{k_{B}T} - 1 \right) + \frac{n_{0}}{p_{0}} \left(\exp\left(\frac{q\phi}{k_{B}T}\right) - \frac{q\phi}{k_{B}T} - 1 \right) \right]$$

At the surface where x = 0, ϕ_s , \mathcal{E}_s

$$\mathcal{E}_{S} = \frac{\sqrt{2}k_{B}T}{qL_{D}}\sqrt{\left(\exp\left(-\frac{q\phi_{S}}{k_{B}T}\right) + \frac{q\phi_{S}}{k_{B}T} - 1\right) + \frac{n_{0}}{p_{0}}\left(\exp\left(\frac{q\phi_{S}}{k_{B}T}\right) - \frac{q\phi_{S}}{k_{B}T} - 1\right)}$$

where
$$L_D = \sqrt{\frac{\varepsilon_S k_B T}{q^2 p_0}}$$
 is the Debye Length





Charge density distribution



Charge density distribution



d = thickness of oxide

Oxide capacitance (unit area)

$$C_i = \frac{c_i}{d}$$

Voltage across oxide
$$V_i = \frac{-Q_s}{C_i} = \frac{-Q_s d}{\varepsilon_i}$$

Applied voltage $V = V_i + \phi_s$

Similar to result for n^+ - p junction

BEFORE strong
inversion
$$W = \sqrt{\frac{2\epsilon_s \phi_s}{qN_A}} \qquad \phi_s < 2\phi_F$$

At strong inversion, depletion region no longer grows, due to screening of interface electrons

strong inversion $\phi_S = 2\phi_F$

$$W_{max} = \sqrt{\frac{2\epsilon_s 2\phi_F}{qN_A}} = 2\sqrt{\frac{\epsilon_s k_B T \ln \frac{N_A}{n_i}}{q^2 N_A}}$$

Depletion charge at threshold

$$W_{max} = \sqrt{\frac{2\epsilon_s 2\phi_F}{qN_A}} \qquad \begin{array}{l} \text{strong inversion limit} \\ \phi_S = 2\phi_F \end{array}$$
$$Q_D = -qN_A W_{max} = -qN_A \sqrt{\frac{2\epsilon_s 2\phi_F}{qN_A}}$$
$$Q_D = -2\sqrt{q\epsilon_s N_A \phi_F}$$

Threshold Voltage (ideal case)



(Assuming that depletion charge dominates Q_s at threshold) Above, $-Q_D$ represents Q_m at threshold, to give the potential drop across the oxide.

Potential drop across the oxide



Electric Field in the oxide is constant in absence of charge

 $Q_m \approx -Q_D$

semiconductor

Voltage drop across oxide

$$\begin{aligned}
C_i^{-1} \text{ (per unit area)} \\
V_i &= \int_0^d \frac{Q_m}{\epsilon_{ox}} \, dx = Q_m \left[\frac{\overline{d}}{\epsilon_{ox}} \right] \approx -\frac{Q_D}{C_i} \\
\text{at threshold}
\end{aligned}$$

$$\epsilon_{ox} = \epsilon_{rox} \epsilon_0$$

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Electric Field distribution



Electric Potential distribution



Summary of conditions – surface potential



Real Surface effects – Work function difference

- We assumed in the previous analysis for simplicity $\Phi_m = \Phi_s$
- In general, we are limited in the choice of metal by technological constraints and

$$\Phi_m \neq \Phi_s$$

It is convenient to define the quantity

$$\Phi_{ms} = \Phi_m - \Phi_s$$

n+ plus polysilicon for gate electrode



Effect of negative workfunction difference



Apply
$$V_{FB} = \Phi_{ms}$$
 to obtain flat band



• Alkali metal ions (e.g. Na^+) inside the oxide cause a mobile charge Q_m inducing negative charge in Si (reduced by careful processing)

• Imperfections in the SiO₂ material cause positive trapped charges Q_{ot} ($\approx 10^{10}$ cm⁻³).

Real Surface effects – Interface charge (2)

- Positive fixed charges Q_f in a transition layer at the interface.
- Positive charges *Q_{it}* at the Si-SiO₂ interface (interface states) due to mismatch causing "ionic" Si atoms with incomplete bonds.

 $Q_{it} + Q_f \approx 10^{10} \text{ cm}^{-3}$ [100] preferred for devices $Q_{it} + Q_f \approx 10^{11} \text{ cm}^{-3}$ [111]

Effect of interface charge



Effect of interface charge



Threshold Voltage



Enhancement and Depletion MOSFET

 Enhancement-mode MOS usually employed for switching elements. These devices are off at zero gate-source voltage.

 Depletion-mode MOS, usually employed to realize "resistors" in logic circuits. These devices are normally on on at zero gate-source voltage.